
On the kernels of Wiener-Hopf-Hankel operators on variable exponent Lebesgue spaces¹

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*Dedicated to honoring the memory of Professor A.V. Balakrishnan,
an eminent researcher and an inspirational communicator.*

Abstract. We investigate properties of the kernels (and cokernels) of Wiener-Hopf plus and minus Hankel operators on variable exponent Lebesgue spaces. Constructive operator identities are used in view to describe those kernels upon the consideration of auxiliary operators. Moreover, a Coburn-Simonenko type theorem is obtained for Wiener-Hopf plus and minus Hankel operators in the framework of variable exponent Lebesgue spaces.

1 Introduction

Convolution type operators [3, 6, 12] arise in a great diversity of applied problems and are useful in a significant amount of different areas of science. Having on their structure a convolution, this can be reflected on the definition of those operators in different ways. Moreover, the convolution itself can be defined in a large amount of varieties, but having always a factorization property associated with a certain integral transform [5, 21, 22, 31]. For the classical convolution case, it is clear that such transform is the Fourier transform, and that the corresponding factorization property ends with the product of the two Fourier transforms of the elements initially used in the convolution. Anyway, now-a-days different convolutions are known and considered in a great diversity of situations and associated with different integral transforms (e.g. like the fractional Fourier transform, Hankel and Hartley type transforms); cf. [4, 5, 21, 22, 23, 31].

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Not rarely, the convolution type operators that arise in the applications appear not isolated but, instead, multiplied by other types of operators. Even if sometimes those other operators are very simple ones, their multiplications by the convolution type operators change significantly the final structural properties of the resulting operators (e.g., their Fredholm and invertibility properties, as well as the structure of their inverses and the corresponding solutions of the equations characterized by them); see [15, 16, 17, 20].

That is the case e.g. in the theory of classical convolution operators firstly defined within some spaces on the real line, which afterwards are truncated to the half-line, and known as Wiener-Hopf operators; or even the previous ones composed with the reflection operator (and therefore giving rise to a Hankel type operator). For the last two types of operators, it also occurs that algebraic sums of those again have special invertibility properties and are also very important in some applications. This is the case in the phenomena of wave diffraction by configurations which exhibit certain symmetries, and where Wiener-Hopf plus Hankel and Wiener-Hopf minus Hankel (i.e. Wiener-Hopf-Hankel) operators appear in a very natural manner.

The classes of Wiener-Hopf plus and minus Hankel operators that initially appear in those wave diffraction problems are considered within a framework of Bessel potential (or Sobolev) spaces; see [7, 8, 9]. Anyway, a lifting procedure is possible to be done in the smoothness of the spaces so that the initial operators are related with new Wiener-Hopf plus and minus Hankel operators acting already between classical Lebesgue spaces.

Having in our disposal the recent huge development in variable exponent Lebesgue spaces [19], a corresponding research has been also started on the analysis of the properties of several types of operators within the framework of variable exponent Lebesgue spaces (cf. [14, 24, 26, 27]), in contrast with the previous knowledge of their properties when considered on the classical Lebesgue space.

In some cases, concrete applications are already known for such more global framework of variable exponents. This is the case e.g. in the modeling of so-called electrorheological fluids [30], thermo-rheological fluids [1] and image processing [13]. In some other cases, the research is just being moved by the natural search for a larger mathematical formal knowledge.

Having all that in mind, in the present paper we will analyse certain invertibility properties of Wiener-Hopf plus Hankel and Wiener-Hopf minus Hankel operators in the framework of variable exponent Lebesgue spaces [10, 11]. Namely, after introducing, in the next section, the formal definition of those operators and spaces, as well as known results of some of their associated operators, in Section 3 we will analyse the kernel of those operators with the help of certain auxiliary operators. To this end, we will use several operator identities between that operators in analysis. In Section 4, we deduce a Coburn-Simonenko type theorem [18] for Wiener-Hopf plus Hankel and Wiener-Hopf minus Hankel operators on variable exponent Lebesgue spaces (which basically describes an alternative result in the dimensions of their kernels and cokernels).

2 Definitions and basic results

Let $p : \mathbb{R} \rightarrow [1, \infty]$ be measurable a.e. finite function. We denote by $L^{p(\cdot)}(\mathbb{R})$ the set of all complex-valued functions f on \mathbb{R} such that

$$I_{p(\cdot)}\left(\frac{f}{\lambda}\right) := \int_{\mathbb{R}} \left| \frac{f(x)}{\lambda} \right|^{p(x)} dx < \infty$$

for some $\lambda > 0$. This set becomes a Banach space when equipped with the norm

$$\|f\|_{p(\cdot)} := \inf \left\{ \lambda > 0 : I_{p(\cdot)} \left(\frac{f}{\lambda} \right) \leq 1 \right\}.$$

The space $L^{p(\cdot)}(\mathbb{R})$ is precisely what we are referring to as the *variable exponent Lebesgue space*, since it generalizes the standard Lebesgue space.

Throughout this paper we will be always assuming that

$$p_- := \operatorname{ess\,inf}_{x \in \mathbb{R}} p(x) > 1, \quad p_+ := \operatorname{ess\,sup}_{x \in \mathbb{R}} p(x) < \infty. \quad (2.1)$$

Under these conditions, the space $L^{p(\cdot)}(\mathbb{R})$ is separable and reflexive, and its dual space is isomorphic to $L^{q(\cdot)}(\mathbb{R})$, where $q(\cdot)$ is the conjugate exponent function defined by

$$\frac{1}{p(x)} + \frac{1}{q(x)} = 1 \quad (x \in \mathbb{R}).$$

Additionally, with condition (2.1) we have that $\|\phi I\|_{\mathcal{L}(L^{p(\cdot)}(\mathbb{R}))} \leq \|\phi\|_{L^\infty(\mathbb{R})}$ for a function $\phi \in L^\infty(\mathbb{R})$, and where I denotes the identity operator.

We shall denote by $L^{p(\cdot)}(\mathbb{R}_+)$ the variable exponent Lebesgue space of complex-valued functions on the positive half-line $\mathbb{R}_+ = (0, +\infty)$. The subspace of $L^{p(\cdot)}(\mathbb{R})$ formed by all functions supported on the closure of \mathbb{R}_+ is denoted by $L_+^{p(\cdot)}(\mathbb{R})$ and $L_-^{p(\cdot)}(\mathbb{R})$ represents the subspace of $L^{p(\cdot)}(\mathbb{R})$ formed by all the functions supported on the closure of $\mathbb{R}_- := (-\infty, 0)$.

We will make an extensive use of the Fourier transformation \mathcal{F} , defined in the Schwartz space $\mathcal{S}(\mathbb{R})$ of rapidly decreasing functions by

$$\mathcal{F} \varphi(\xi) := \int_{\mathbb{R}} e^{i\eta\xi} \varphi(\eta) d\eta \quad (\xi \in \mathbb{R}),$$

and which has an inverse \mathcal{F}^{-1} , also given on $\mathcal{S}(\mathbb{R})$ by

$$\mathcal{F}^{-1} \psi(\eta) := \frac{1}{2\pi} \int_{\mathbb{R}} e^{-i\eta\xi} \psi(\xi) d\xi \quad (\eta \in \mathbb{R}).$$

Definition 1. (cf. [25, 29]) *If $1 < p_- \leq p_+ < \infty$, a function $\phi \in L^\infty(\mathbb{R})$ is said to be an $L^{p(\cdot)}$ -Fourier multiplier if there is a constant C such that for all $f \in \mathcal{S}(\mathbb{R})$ we have*

$$\|\mathcal{F}^{-1} \phi \cdot \mathcal{F} f\|_{p(\cdot)} \leq C \|f\|_{p(\cdot)}.$$

For any $L^{p(\cdot)}$ -Fourier multiplier ϕ the operator $f \mapsto \mathcal{F}^{-1} \phi \cdot \mathcal{F} f$ extends uniquely to a bounded operator on $L^{p(\cdot)}(\mathbb{R})$ which will be denoted by W_ϕ^0 . The set of all $L^{p(\cdot)}$ -Fourier multipliers will be denoted by $\mathcal{M}_{p(\cdot)}$. It is clear that $\mathcal{M}_{p(\cdot)}$ is a unital normed algebra under pointwise operations and the norm

$$\|\phi\|_{\mathcal{M}_{p(\cdot)}} := \|W_\phi^0\|_{\mathcal{L}(L^{p(\cdot)}(\mathbb{R}))} = \|\mathcal{F}^{-1} \phi \cdot \mathcal{F}\|_{\mathcal{L}(L^{p(\cdot)}(\mathbb{R}))}.$$

We are now in condition to identify in a mathematical way the main objects of this work.

We will consider Wiener-Hopf plus and minus Hankel operators, acting between Lebesgue spaces with variable exponent $p(\cdot)$, denoted by

$$W_\phi \pm H_\phi : L_+^{p(\cdot)}(\mathbb{R}) \rightarrow L^{p(\cdot)}(\mathbb{R}_+),$$

with W_ϕ and H_ϕ being Wiener-Hopf and Hankel operators defined by

$$W_\phi = r_+ \mathcal{F}^{-1} \phi \cdot \mathcal{F}, \quad H_\phi = r_+ \mathcal{F}^{-1} \phi \cdot \mathcal{F} J, \quad (2.2)$$

respectively. Here, r_+ represents the operator of restriction from $L^{p(\cdot)}(\mathbb{R})$ onto $L^{p(\cdot)}(\mathbb{R}_+)$, ϕ is the so-called Fourier symbol which is assumed to belong to $\mathcal{M}_{p(\cdot)}$, and

$$J : L_+^{p(\cdot)}(\mathbb{R}) \rightarrow L^{p(\cdot)}(\mathbb{R})$$

is the reflection operator given by the rule

$$J\varphi(x) = \tilde{\varphi}(x) = \varphi(-x)$$

which throughout the paper will be always defined for even functions $p(\cdot)$ (so that J will therefore be a bounded operator in those variable exponent Lebesgue spaces).

Under those conditions, the Wiener-Hopf and the Hankel operators, defined in (2.2), are bounded linear operators.

It is also interesting to remark that the boundedness of a wide variety of operators follows from the boundedness of the Hardy-Littlewood maximal operator on variable exponent Lebesgue spaces. Given $f \in L_{\text{loc}}^1(\mathbb{R})$, we recall that the Hardy-Littlewood maximal operator M is defined by

$$(Mf)(x) := \sup_{x \in \Omega} \frac{1}{|\Omega|} \int_{\Omega} |f(y)| dy,$$

where the supremum is taken over all intervals $\Omega \subset \mathbb{R}$ containing x , and the Cauchy singular integral operator S is defined by

$$(Sf)(x) := \frac{1}{\pi i} \int_{\mathbb{R}} \frac{f(\tau)}{\tau - x} d\tau,$$

where the integral is understood in the principal value sense.

Theorem 1. (cf. e.g., [26, Theorem 2.1.]) Let $p : \mathbb{R} \rightarrow [1, \infty]$ be a measurable function satisfying (2.1). If the Hardy-Littlewood maximal operator M is bounded on $L^{p(\cdot)}(\mathbb{R})$, then the Cauchy singular integral operator S is bounded on $L^{p(\cdot)}(\mathbb{R})$.

The following result states a sufficient condition on $p(\cdot)$ for M to be bounded on $L^{p(\cdot)}(\mathbb{R})$.

Theorem 2. (cf. e.g. [29, Theorem 2.5.], [19]) Let $p : \mathbb{R} \rightarrow [1, \infty]$ satisfy (2.1). In addition, suppose that there exist constants A_0 and A_∞ such that $p(\cdot)$ satisfies

$$|p(x) - p(y)| \leq \frac{A_0}{-\log|x-y|}, \quad |x-y| \leq \frac{1}{2}, \quad (2.3)$$

and

$$|p(x) - p(y)| \leq \frac{A_\infty}{\log(e+|x|)}, \quad |x| \leq |y|. \quad (2.4)$$

Then, the Hardy-Littlewood maximal operator is bounded on $L^{p(\cdot)}(\mathbb{R})$.

It is important to stress that the log-Hölder continuity conditions (2.3)-(2.4) are not necessary conditions for the last purpose. The exponent $p(\cdot)$ can even be discontinuous. In fact, Lerner have shown in [28] that there exist discontinuous variable exponents $p(\cdot)$ (without limit at infinity) for which the Hardy-Littlewood maximal operator is bounded on the $L^{p(\cdot)}(\mathbb{R})$ space. Under conditions

(2.1), (2.3) and (2.4), from [25] it is also known that $\mathcal{M}_{p(\cdot)}$ is a Banach algebra under pointwise operations and the norm $\|\cdot\|_{\mathcal{M}_{p(\cdot)}}$, and that $\mathcal{M}_{p(\cdot)}$ is continuously embedded into $L^\infty(\mathbb{R})$.

Let $\mathcal{B}(\mathbb{R})$ denote the class of exponents $p : \mathbb{R} \rightarrow [1, \infty]$ continuous on \mathbb{R} satisfying (2.1), (2.3) and (2.4) (and therefore with M and S being bounded operators on $L^{p(\cdot)}(\mathbb{R})$). Additionally, $\mathcal{B}_e(\mathbb{R})$ represents the set of all even functions $p(\cdot) \in \mathcal{B}(\mathbb{R})$.

Let ℓ_0 denote the zero extension operator from the space $L^{p(\cdot)}(\mathbb{R}_+)$ into the space $L_+^{p(\cdot)}(\mathbb{R})$,

$$\ell_0 : L^{p(\cdot)}(\mathbb{R}_+) \rightarrow L_+^{p(\cdot)}(\mathbb{R}). \quad (2.5)$$

We will denote by P the canonical projection of $L^{p(\cdot)}(\mathbb{R})$ onto $L_+^{p(\cdot)}(\mathbb{R})$, and its complementary projection $Q := I - P$ of $L^{p(\cdot)}(\mathbb{R})$ onto $L_-^{p(\cdot)}(\mathbb{R})$. We have that

$$P := \ell_0 r_+.$$

Note that $P^2 = P$ and $Q^2 = Q$.

Additionally, we will use the notation \mathcal{P} to denote the Riesz projection defined by

$$\mathcal{P} = \mathcal{F} \ell_0 r_+ \mathcal{F}^{-1}$$

and denote by Q its complementary projection, $Q = I - \mathcal{P}$.

By analogy with [3, Proposition 2.10], we derive the following relations between Wiener-Hopf and Hankel operators acting on variable exponent Lebesgue spaces.

Proposition 1. Let $\phi, \varphi \in \mathcal{M}_{p(\cdot)}$ and $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. Then

$$W_{\phi\varphi} = W_\phi \ell_0 W_\varphi + H_\phi \ell_0 H_{\tilde{\varphi}} \quad (2.6)$$

$$H_{\phi\varphi} = W_\phi \ell_0 H_\varphi + H_\phi \ell_0 W_{\tilde{\varphi}}. \quad (2.7)$$

In what follows, we will also make use of the relations

$$JQ = PJ, \quad JP = QJ, \quad J^2 = I, \quad JW_\phi^0 J = W_\phi^0.$$

3 Kernels of Wiener-Hopf-Hankel operators

We will denote by $\mathcal{G}\mathcal{M}_{p(\cdot)}$ the group of all invertible elements of $\mathcal{M}_{p(\cdot)}$. Let $\phi \in \mathcal{G}\mathcal{M}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$ and consider the Wiener-Hopf plus and minus Hankel operators,

$$W_\phi \pm H_\phi : L_+^{p(\cdot)}(\mathbb{R}) \rightarrow L^{p(\cdot)}(\mathbb{R}_+).$$

With some computations, we can obtain an important formula to “associate” these operators with a (pure) Wiener-Hopf operator. Namely, multiplying $W_\phi \pm H_\phi$ on the left by the zero extension operator $\ell_0 : L^{p(\cdot)}(\mathbb{R}_+) \rightarrow L_+^{p(\cdot)}(\mathbb{R})$, and on the right by the projection $P : L^{p(\cdot)}(\mathbb{R}) \rightarrow L_+^{p(\cdot)}(\mathbb{R})$, we obtain

$$\ell_0(W_\phi \pm H_\phi)P = P(W_\phi^0 \pm W_\phi^0 J)P : L^{p(\cdot)}(\mathbb{R}) \rightarrow L_+^{p(\cdot)}(\mathbb{R}),$$

where we recall that $W_\phi^0 = \mathcal{F}^{-1}\phi \cdot \mathcal{F}$.

To extend these operators to the full $L^{p(\cdot)}(\mathbb{R})$ space, we use the projection $Q : L^{p(\cdot)}(\mathbb{R}) \rightarrow L_-^{p(\cdot)}(\mathbb{R})$, obtaining the equivalent after extension [2] operators

$$PW_\phi^0 P \pm PW_\phi^0 J P + Q : L^{p(\cdot)}(\mathbb{R}) \rightarrow L^{p(\cdot)}(\mathbb{R}).$$

Again, let $I := I_{L^{p(\cdot)}(\mathbb{R})}$ denote the identity operator on $L^{p(\cdot)}(\mathbb{R})$. Observe that

$$PW_\phi^0 P \pm PW_\phi^0 J P + Q = (I - PW_\phi^0 Q \mp PW_\phi^0 J Q)(PW_\phi^0 \pm PW_\phi^0 J + Q),$$

where the operators $I - PW_\phi^0 Q \mp PW_\phi^0 J Q$ are invertible (with inverses given by $I + PW_\phi^0 Q \pm PW_\phi^0 J Q$, respectively). Let

$$\Psi = \begin{bmatrix} 0 & \widetilde{\phi\phi^{-1}} \\ -1 & \widetilde{\phi^{-1}} \end{bmatrix}. \quad (3.1)$$

We can state the following

$$\begin{bmatrix} PW_\phi^0 P + PW_\phi^0 J P + Q & 0 \\ 0 & PW_\phi^0 P - PW_\phi^0 J P + Q \end{bmatrix} = A (PW_\Psi^0 P + Q) B \quad (3.2)$$

where $A := A_1 A_2$ and B are invertible operators with

$$\begin{aligned} A_1 &= \frac{1}{2} \begin{bmatrix} I - PW_\phi^0 Q - PW_\phi^0 J Q & 0 \\ 0 & I - PW_\phi^0 Q + PW_\phi^0 J Q \end{bmatrix} \begin{bmatrix} I & J \\ I & -J \end{bmatrix} \\ A_2 &= I + PW_\Psi^0 Q \\ B &= \begin{bmatrix} I & 0 \\ W_\phi^0 & W_\phi^0 \end{bmatrix} \begin{bmatrix} I & I \\ J & -J \end{bmatrix}, \end{aligned}$$

with inverses given by

$$\begin{aligned} A_1^{-1} &= \begin{bmatrix} I & I \\ J & -J \end{bmatrix} \begin{bmatrix} I + PW_\phi^0 Q + PW_\phi^0 J Q & 0 \\ 0 & I + PW_\phi^0 Q - PW_\phi^0 J Q \end{bmatrix} \\ A_2^{-1} &= I - PW_\Psi^0 Q \\ B^{-1} &= \frac{1}{2} \begin{bmatrix} I & J \\ I & -J \end{bmatrix} \begin{bmatrix} I & 0 \\ -I & W_{\phi^{-1}}^0 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} I - J & W_{\phi^{-1}}^0 J \\ I + J & -W_{\phi^{-1}}^0 J \end{bmatrix}, \end{aligned}$$

respectively.

Moreover, we can rewrite the operator $PW_\Psi^0 P + Q$ as a product of two matrix operators. Namely, for the first component, we have

$$PW_\Psi^0 P = \begin{bmatrix} PW_{\phi\phi^{-1}}^0 P & 0 \\ 0 & I_{L_+^{p(\cdot)}(\mathbb{R})} \end{bmatrix} \begin{bmatrix} 0 & I_{L_+^{p(\cdot)}(\mathbb{R})} \\ -I_{L_+^{p(\cdot)}(\mathbb{R})} & PW_{\phi^{-1}}^0 P \end{bmatrix}, \quad (3.3)$$

where

$$C := \begin{bmatrix} 0 & I_{L_+^{p(\cdot)}(\mathbb{R})} \\ -I_{L_+^{p(\cdot)}(\mathbb{R})} & PW_{\phi^{-1}}^0 P \end{bmatrix}$$

is invertible with inverse

$$C^{-1} = \begin{bmatrix} PW_{\tilde{\phi}^{-1}}^0 P - I_{L_+^{p(\cdot)}(\mathbb{R})} & \\ I_{L_+^{p(\cdot)}(\mathbb{R})} & 0 \end{bmatrix}.$$

Thus, we can rewrite (3.2) as

$$\begin{bmatrix} PW_{\tilde{\phi}}^0 P + PW_{\tilde{\phi}}^0 J P + Q & 0 \\ 0 & PW_{\tilde{\phi}}^0 P - PW_{\tilde{\phi}}^0 J P + Q \end{bmatrix} = A \left(\begin{bmatrix} PW_{\tilde{\phi}^{-1}}^0 P & 0 \\ 0 & I_{L_+^{p(\cdot)}(\mathbb{R})} \end{bmatrix} C + Q \right) B. \quad (3.4)$$

Theorem 3. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_c(\mathbb{R})$ and consider the operators

$$W_{\phi} \pm H_{\phi} : L_+^{p(\cdot)}(\mathbb{R}) \rightarrow L^{p(\cdot)}(\mathbb{R}_+).$$

(i) If $f \in \ker W_{\tilde{\phi}^{-1}}$, then

$$\begin{aligned} (F, G)^T &:= \frac{1}{2} \left(PW_{\tilde{\phi}^{-1}}^0 f + JQW_{\tilde{\phi}^{-1}}^0 f, PW_{\tilde{\phi}^{-1}}^0 f - JQW_{\tilde{\phi}^{-1}}^0 f \right)^T \\ &\in \ker \operatorname{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}); \end{aligned} \quad (3.5)$$

(ii) If $(F, G)^T \in \ker \operatorname{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi})$, then

$$f := PW_{\tilde{\phi}}^0 (F + G) + PW_{\tilde{\phi}}^0 J (F - G) \in \ker W_{\tilde{\phi}^{-1}}. \quad (3.6)$$

Moreover, the operators

$$\begin{aligned} L_1 &: \ker W_{\tilde{\phi}^{-1}} \rightarrow \ker \operatorname{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}), \\ L_2 &: \ker \operatorname{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}) \rightarrow \ker W_{\tilde{\phi}^{-1}}, \end{aligned}$$

defined by the actions in (3.5) and (3.6) are invertible operators satisfying $L_2 = L_1^{-1}$.

Proof. The result follows from relation (3.4). In fact, from (3.4), if $f \in \ker W_{\tilde{\phi}^{-1}}$, and consequently

$$(f, 0)^T \in \ker \operatorname{diag}(W_{\tilde{\phi}^{-1}}, I),$$

then $B^{-1}PC^{-1}((f, 0)^T)$ belongs to

$$\ker(W_{\phi} + H_{\phi} + Q, W_{\phi} - H_{\phi} + Q) \cong \ker(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}).$$

Moreover,

$$B^{-1}PC^{-1}((f, 0)^T) = \frac{1}{2} \left(PW_{\tilde{\phi}^{-1}}^0 P f + JQW_{\tilde{\phi}^{-1}}^0 P f, PW_{\tilde{\phi}^{-1}}^0 P f - JQW_{\tilde{\phi}^{-1}}^0 P f \right)^T,$$

which is equal to $\frac{1}{2} \left(PW_{\tilde{\phi}^{-1}}^0 f + JQW_{\tilde{\phi}^{-1}}^0 f, PW_{\tilde{\phi}^{-1}}^0 f - JQW_{\tilde{\phi}^{-1}}^0 f \right)^T$ since $f \in L_+^{p(\cdot)}(\mathbb{R})$. Thus, we obtain (3.5).

On the other hand, if $(F, G)^T \in \ker \operatorname{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi})$, then

$$CPB((F, G)^T) \in \ker \operatorname{diag}(W_{\tilde{\phi}^{-1}}, I).$$

Thus,

$$CPB((F, G)^T) = \begin{bmatrix} PW_{\tilde{\phi}}^0(F + G) + PW_{\tilde{\phi}}^0 J(F - G) \\ -(F + G) + PW_{\tilde{\phi}-1}^0 PW_{\tilde{\phi}}^0(F + G) + PW_{\tilde{\phi}-1}^0 PW_{\tilde{\phi}}^0 J(F - G) \end{bmatrix} \in \ker \text{diag}(W_{\tilde{\phi}\tilde{\phi}-1}, I),$$

which proves (3.6).

Finally, let f and $(F, G)^T$ as defined in (3.5) and (3.6). We observe that

$$CPB(F, G)^T = (f, 0)^T,$$

which completes the proof.

Proposition 2. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. If $f \in \ker W_{\tilde{\phi}\tilde{\phi}-1}$, then

$$JQW_{\tilde{\phi}\tilde{\phi}-1}^0 f \in \ker W_{\tilde{\phi}\tilde{\phi}-1}$$

and $(JQW_{\tilde{\phi}\tilde{\phi}-1}^0)^2 f = f$.

Proof. Let $\phi \in \mathcal{GM}_{p(\cdot)}$ ($p(\cdot) \in \mathcal{B}_e(\mathbb{R})$) and $f \in \ker W_{\tilde{\phi}\tilde{\phi}-1}$. Then

$$\begin{aligned} \ell_0 W_{\tilde{\phi}\tilde{\phi}-1} (JQW_{\tilde{\phi}\tilde{\phi}-1}^0 f) &= PW_{\tilde{\phi}\tilde{\phi}-1}^0 PJQW_{\tilde{\phi}\tilde{\phi}-1}^0 f = JQW_{\tilde{\phi}\tilde{\phi}-1}^0 QW_{\tilde{\phi}\tilde{\phi}-1}^0 f \\ &= JQW_{\tilde{\phi}\tilde{\phi}-1}^0 W_{\tilde{\phi}\tilde{\phi}-1}^0 f - JQW_{\tilde{\phi}\tilde{\phi}-1}^0 PW_{\tilde{\phi}\tilde{\phi}-1}^0 f \\ &= JQf - JQW_{\tilde{\phi}\tilde{\phi}-1}^0 0 \\ &= 0. \end{aligned}$$

Additionally, for any $f \in \ker W_{\tilde{\phi}\tilde{\phi}-1}$, we have

$$\begin{aligned} (JQW_{\tilde{\phi}\tilde{\phi}-1}^0)^2 f &= JQW_{\tilde{\phi}\tilde{\phi}-1}^0 PJQW_{\tilde{\phi}\tilde{\phi}-1}^0 Pf = PW_{\tilde{\phi}\tilde{\phi}-1}^0 QW_{\tilde{\phi}\tilde{\phi}-1}^0 f \\ &= PW_{\tilde{\phi}\tilde{\phi}-1}^0 W_{\tilde{\phi}\tilde{\phi}-1}^0 f - PW_{\tilde{\phi}\tilde{\phi}-1}^0 PW_{\tilde{\phi}\tilde{\phi}-1}^0 f = f, \end{aligned}$$

which completes the proof.

Consider the operator $P_{\tilde{\phi}\tilde{\phi}-1} := JQW_{\tilde{\phi}\tilde{\phi}-1}^0 P|_{\ker W_{\tilde{\phi}\tilde{\phi}-1}}$. From Proposition 2,

$$P_{\tilde{\phi}\tilde{\phi}-1} : \ker W_{\tilde{\phi}\tilde{\phi}-1} \rightarrow \ker W_{\tilde{\phi}\tilde{\phi}-1}$$

and

$$P_{\tilde{\phi}\tilde{\phi}-1}^2 = I.$$

Additionally, $P_{\tilde{\phi}\tilde{\phi}-1}^- := (1/2)(I - P_{\tilde{\phi}\tilde{\phi}-1})$ and $P_{\tilde{\phi}\tilde{\phi}-1}^+ := (1/2)(I + P_{\tilde{\phi}\tilde{\phi}-1})$ are complementary projections generating a decomposition of $\ker W_{\tilde{\phi}\tilde{\phi}-1}$.

For $f \in \ker W_{\tilde{\phi}\tilde{\phi}-1}$, let

$$\varphi^\pm(f) := \frac{1}{2}(PW_{\tilde{\phi}\tilde{\phi}-1}^0 f \pm JQW_{\tilde{\phi}\tilde{\phi}-1}^0 f).$$

By Theorem 3, $\varphi^\pm(f) \in \ker(W_\phi \pm H_\phi)$.

Lemma 1. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. For every $f \in \ker W_{\phi\tilde{\phi}^{-1}}$, the following relations

$$(W_{\tilde{\phi}} + H_{\tilde{\phi}})\varphi_+(f) = P_{\phi\tilde{\phi}^{-1}}^+ f, \quad (W_{\tilde{\phi}} - H_{\tilde{\phi}})\varphi_-(f) = P_{\phi\tilde{\phi}^{-1}}^- f$$

hold. The corresponding mapping $\varphi_+ : \text{im } P_{\phi\tilde{\phi}^{-1}}^+ \rightarrow \text{im } P_{\phi\tilde{\phi}^{-1}}^+$ and $\varphi_- : \text{im } P_{\phi\tilde{\phi}^{-1}}^- \rightarrow \text{im } P_{\phi\tilde{\phi}^{-1}}^-$ are injective operators.

Proof. Let $f \in \ker W_{\phi\tilde{\phi}^{-1}}$. We will show that the operator $W_{\tilde{\phi}} + H_{\tilde{\phi}}$ maps $\varphi^+(f)$ into $P_{\phi\tilde{\phi}^{-1}}^+ f$ and the operator $W_{\tilde{\phi}} - H_{\tilde{\phi}}$ maps $\varphi^-(f)$ into $P_{\phi\tilde{\phi}^{-1}}^- f$.

For the first case, using Proposition 1, we have

$$\begin{aligned} \ell_0(W_{\tilde{\phi}} + H_{\tilde{\phi}})\varphi^+(f) &= \frac{1}{2}\ell_0(W_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + W_{\tilde{\phi}}\ell_0JQW_{\tilde{\phi}^{-1}}^0f + H_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + H_{\tilde{\phi}}\ell_0JQW_{\tilde{\phi}^{-1}}^0f) \\ &= \frac{1}{2}(\ell_0W_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + \ell_0W_{\tilde{\phi}}PW_{\tilde{\phi}^{-1}}^0Jf + \ell_0H_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + \ell_0H_{\tilde{\phi}}PW_{\tilde{\phi}^{-1}}^0Jf) \\ &= \frac{1}{2}(\ell_0W_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + \ell_0W_{\tilde{\phi}}\ell_0H_{\phi^{-1}}f + \ell_0H_{\tilde{\phi}}\ell_0W_{\tilde{\phi}^{-1}}f + \ell_0H_{\tilde{\phi}}\ell_0H_{\phi^{-1}}f) \\ &= \frac{1}{2}(\ell_0W_{\phi\tilde{\phi}^{-1}}f + \ell_0H_{\phi\tilde{\phi}^{-1}}f) \\ &= \frac{1}{2}(f + PW_{\phi\tilde{\phi}^{-1}}^0Jf) \\ &= \frac{1}{2}(f + PJW_{\phi\tilde{\phi}^{-1}}^0f) \\ &= \frac{1}{2}(f + JQW_{\phi\tilde{\phi}^{-1}}^0f) = P_{\phi\tilde{\phi}^{-1}}^+ f. \end{aligned}$$

The other case is proved analogously.

Proposition 3. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. Then

$$\begin{aligned} \ker(W_{\phi} + H_{\phi}) &= \varphi^+(\text{im } P_{\phi\tilde{\phi}^{-1}}^+), \\ \ker(W_{\phi} - H_{\phi}) &= \varphi^-(\text{im } P_{\phi\tilde{\phi}^{-1}}^-) \end{aligned}$$

Proof. Using the invertible operator L_1 defined in Theorem 3, we obtain

$$\ker \text{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}) = L_1(\ker W_{\phi\tilde{\phi}^{-1}}),$$

where

$$\ker W_{\phi\tilde{\phi}^{-1}} = \text{im } P_{\phi\tilde{\phi}^{-1}}^+ \dot{+} \text{im } P_{\phi\tilde{\phi}^{-1}}^-.$$

Hence,

$$\ker \text{diag}(W_{\phi} + H_{\phi}, W_{\phi} - H_{\phi}) = L_1(\ker W_{\phi\tilde{\phi}^{-1}}) = L_1(\text{im } P_{\phi\tilde{\phi}^{-1}}^+) \dot{+} L_1(\text{im } P_{\phi\tilde{\phi}^{-1}}^-).$$

Thus, having in mind the definitions of L_1 , L_2 , φ^+ and φ^- , the result follows.

Notice that from the equivalence relation (3.4) we have that if $W_{\phi\tilde{\phi}^{-1}}$ is a Fredholm operator (and continuing to assume that $\phi \in \mathcal{GM}_{p(\cdot)}$ and $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$), then both operators $W_{\phi} + H_{\phi}$ and $W_{\phi} - H_{\phi}$ are Fredholm and it holds the following identity for the Fredholm indices of the related operators:

$$\text{ind}(W_{\phi} + H_{\phi}) + \text{ind}(W_{\phi} - H_{\phi}) = \text{ind}W_{\phi\tilde{\phi}^{-1}}.$$

In more detail, and now from Proposition 3, we immediately have the following result for the dimensions of the kernels of $W_{\phi} + H_{\phi}$ and $W_{\phi} - H_{\phi}$.

Corollary 1. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. Then

$$\begin{aligned}\dim \ker(W_\phi + H_\phi) &= \dim \operatorname{im} P_{\phi\tilde{\phi}^{-1}}^+, \\ \dim \ker(W_\phi - H_\phi) &= \dim \operatorname{im} P_{\phi\tilde{\phi}^{-1}}^-.\end{aligned}$$

4 Coburn-Simonenko type theorem for Wiener-Hopf-Hankel operators

The following theorem is the analogue of Coburn-Simonenko Theorem for Wiener-Hopf operators on variable exponent Lebesgue spaces, which states that a nonzero bounded Wiener-Hopf operator has a trivial kernel or a dense range.

Theorem 4. If $\phi \in \mathcal{M}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}(\mathbb{R})$, and if ϕ does not vanish identically, then the kernel of W_ϕ in $L_+^{p(\cdot)}(\mathbb{R})$ is trivial or the image of W_ϕ is dense in $L_+^{p(\cdot)}(\mathbb{R})$.

Proof. First, recall that W_ϕ is equivalent after extension [2] to $\phi\mathcal{P} + Q$.

We will show that $\phi\mathcal{P} + Q$ is injective on $L^{p(\cdot)}(\mathbb{R})$ or $Q\phi I + \mathcal{P}$ is injective on $L^{q(\cdot)}(\mathbb{R})$.

Assume the contrary, i.e., assume that there are nonzero $f \in L^{p(\cdot)}(\mathbb{R})$ and $g \in L^{q(\cdot)}(\mathbb{R})$ such that

$$(\phi\mathcal{P} + Q)f = 0 \quad \text{and} \quad (Q\phi I + \mathcal{P})g = 0.$$

The first identity implies that

$$f_- := Qf \in QL^{p(\cdot)}(\mathbb{R}), \quad f_+ := \mathcal{P}f \in \mathcal{P}L^{p(\cdot)}(\mathbb{R})$$

and

$$\phi f_+ + f_- = 0.$$

From the second equality, we have that $Q\phi g = 0$ and $\mathcal{P}g = 0$, whence $g_+ := \phi g \in \mathcal{P}L^{q(\cdot)}(\mathbb{R})$ and $g_- := g \in QL^{q(\cdot)}(\mathbb{R})$. Multiplying the equality $\phi f_+ = -f_-$ by g we get $\phi g f_+ = -f_- g_-$ and hence,

$$g_+ f_+ = -f_- g_-. \quad (4.1)$$

The left-hand side of (4.1) belongs to $\mathcal{P}L^1(\mathbb{R})$ while its right-hand side belongs to $QL^1(\mathbb{R})$. Consequently, $g_+ f_+ = 0$ and $f_- g_- = 0$. Since $g_- = g \neq 0$, then $f_- = 0$. It follows that $f_+ \neq 0$, otherwise, $f = f_- + f_+ = 0$.

Now, the equality $g_+ f_+ = 0$ shows that $0 = g_+ = \phi g$. Since $\phi \in \mathcal{M}_{p(\cdot)} \setminus \{0\}$, we arrive at conclusion that $g = g_- = 0$, which contradicts the hypothesis.

If $\phi\mathcal{P} + Q$ is injective on $L^{p(\cdot)}(\mathbb{R})$, then W_ϕ is injective on $L^{p(\cdot)}(\mathbb{R})$. On the other hand, if $Q\phi I + \mathcal{P}$ is injective on $L^{q(\cdot)}(\mathbb{R})$, then $C(Q\phi I + \mathcal{P})C$ is also injective on $L^{p(\cdot)}(\mathbb{R})$ (where $C : \varphi \mapsto \bar{\varphi}$ denote the operator of complex conjugation in $L^{p(\cdot)}(\mathbb{R})$). Since $C(Q\phi I + \mathcal{P})C = (\phi\mathcal{P} + Q)^*$, we infer that $\phi\mathcal{P} + Q$ has a dense range on $L^{p(\cdot)}$ and thus, W_ϕ has a dense range on $L^{p(\cdot)}(\mathbb{R}_+)$.

Theorem 5. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$, and let T denote any of the operators $W_\phi - H_\phi$, $W_\phi + H_\phi$. Then, at least, one of the spaces $\ker T$ or $\operatorname{coker} T$ is trivial.

Proof. If $\dim \ker W_{\phi\tilde{\phi}^{-1}} > 0$, then the Coburn-Simonenko theorem for Wiener-Hopf operators gives that

$$\operatorname{coker} W_{\phi\tilde{\phi}^{-1}} = \{0\}.$$

From identity (3.4), we obtain that $\text{coker}(W_\phi \pm H_\phi) = \{0\}$.

On the other hand, assume that $\ker W_{\phi_{\tilde{\phi}^{-1}}} = \{0\}$. Then, the equivalence after extension relation (3.4) implies that

$$\ker(W_\phi \pm H_\phi) = \{0\}.$$

Thus, we conclude that the Coburn-Simonenko theorem is valid for Wiener-Hopf plus Hankel, and Wiener-Hopf minus Hankel operators acting on variable exponent Lebesgue spaces.

Corollary 2. Let $\phi \in \mathcal{GM}_{p(\cdot)}$, $p(\cdot) \in \mathcal{B}_e(\mathbb{R})$. If $\dim \ker W_{\phi_{\tilde{\phi}^{-1}}} = 0$, then $\ker(W_\phi \pm H_\phi) = \{0\}$ and if $\dim \ker W_{\phi_{\tilde{\phi}^{-1}}} > 0$, then $\text{coker}(W_\phi \pm H_\phi) = \{0\}$.

Proof. This statement is directly derived from the proof of the last result.

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